

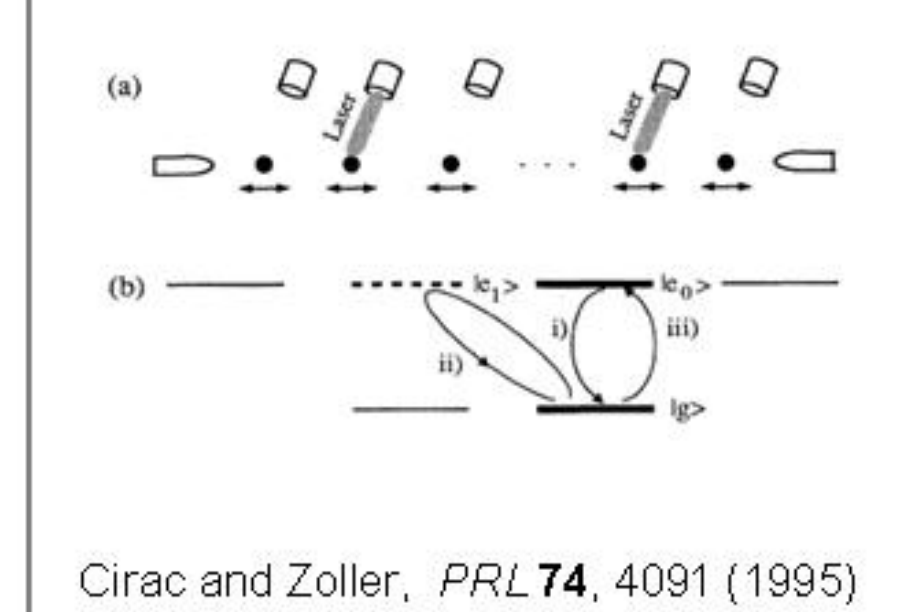
Networking surface electrode ion traps for large-scale QIP

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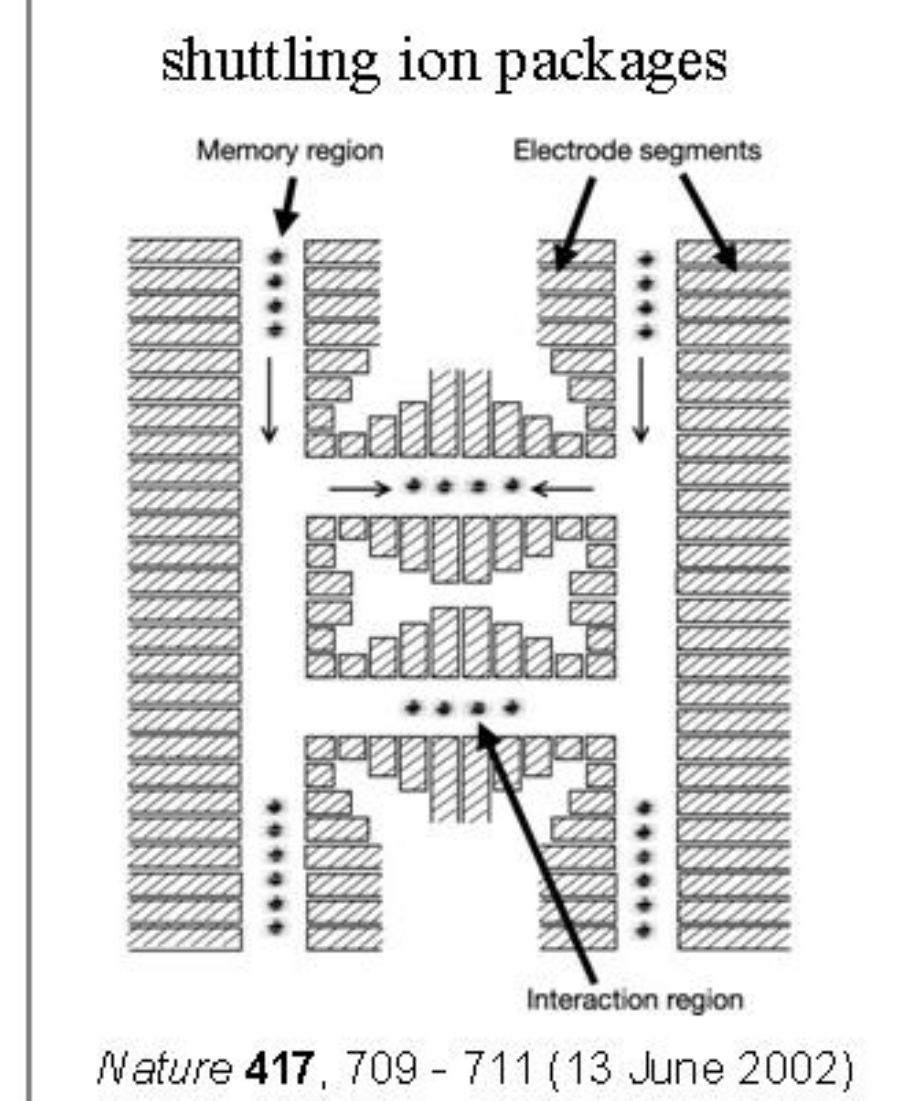
Towards an Ion Trap Quantum Processor

original proposal: ions in a string



Cirac and Zoller, *PRL* **74**, 4091 (1995)

improved proposal:



Nature **417**, 709 - 711 (13 June 2002)

- Trapped ions constitute one of the most promising systems to implement scalable quantum computation

- all DiVincenzo criteria demonstrated, but
How can we scale up?

- in a single trap scaling was found to be impossible because ions in a long string require an extremely accurate control of oscillation frequencies (also extremely large radial frequencies), upper left figure

- another proposal (bottom left figure) suggested to use ion packages and to transport quantum information by shuttling them through microstructures

- 2D is required to separate and join arbitrary groups of ions (e.g. necessary for entangling gates)

- sophisticated layouts implement memory units & interaction zones

How does this look like?

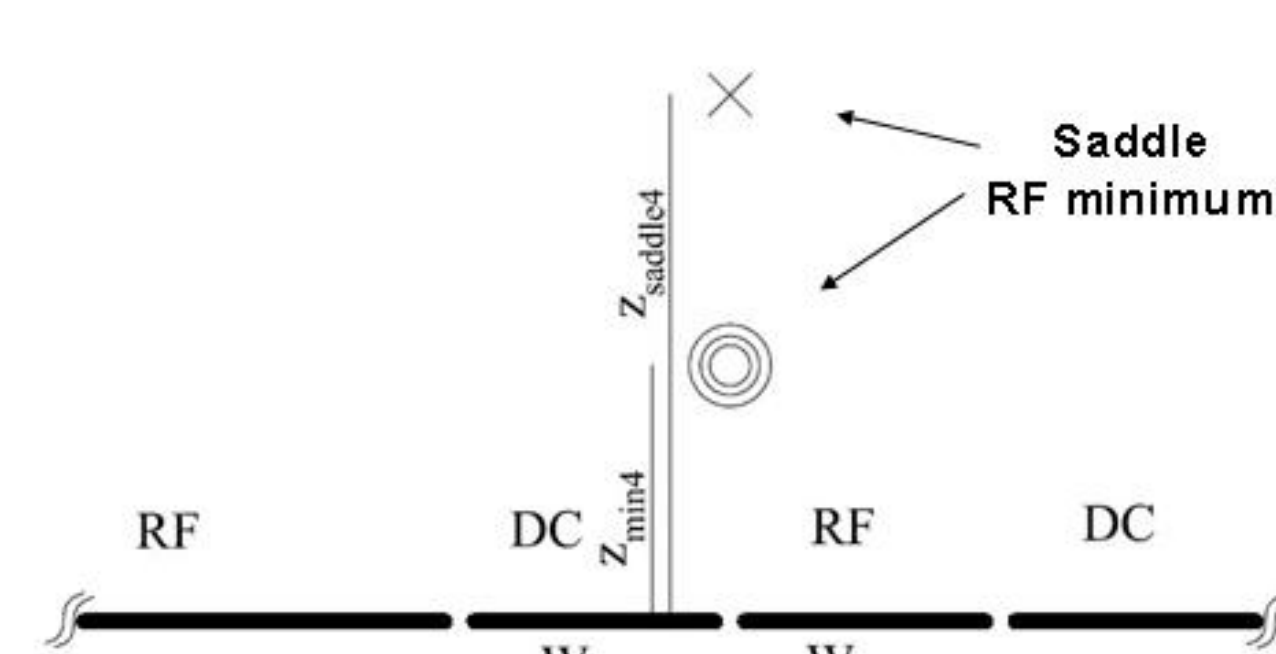
problematic issues

- junctions have RF barriers
- technological scalable trap layouts (simplest approach: all in a single plane)
- required resources still enormous
- ..., probably still more

General properties of surface traps: Cross Section Model (2D)

- we solve the Laplace equation with Dirichlet boundary conditions analytically and use some idealizations from which we can derive general properties of the RF potentials of surface traps and tabulate them.
- wires extend in the 3rd dimension (out of the paper plane) and are assumed infinitely long

'four wire surface traps' /
(asymmetric four wire layouts+ground)



Position of RF minimum and saddle

$$z_{\min 4} = W$$

$$z_{\text{saddle}4} = \sqrt{2 + \sqrt{5}} W \sim 2.058W,$$

Trapping depth (t.d.)

$$t.d.4(W) = \frac{Q^2 V_{RF}^2}{4m\Omega^2 z_{\min 4}^2} \cdot \frac{2}{\pi^2 (11 + 5\sqrt{5})}$$

$$\equiv t.d.\text{-hyp} \sim 0.00914$$

Trapping Frequencies

$$\omega_z \equiv \omega_x = \frac{QV_{RF}}{\sqrt{2}m\Omega z_{\min 4}^2} \cdot \frac{1}{\pi}$$

$$\equiv \omega_{\text{hyp}} \sim 0.318 \text{ reduction factor}$$

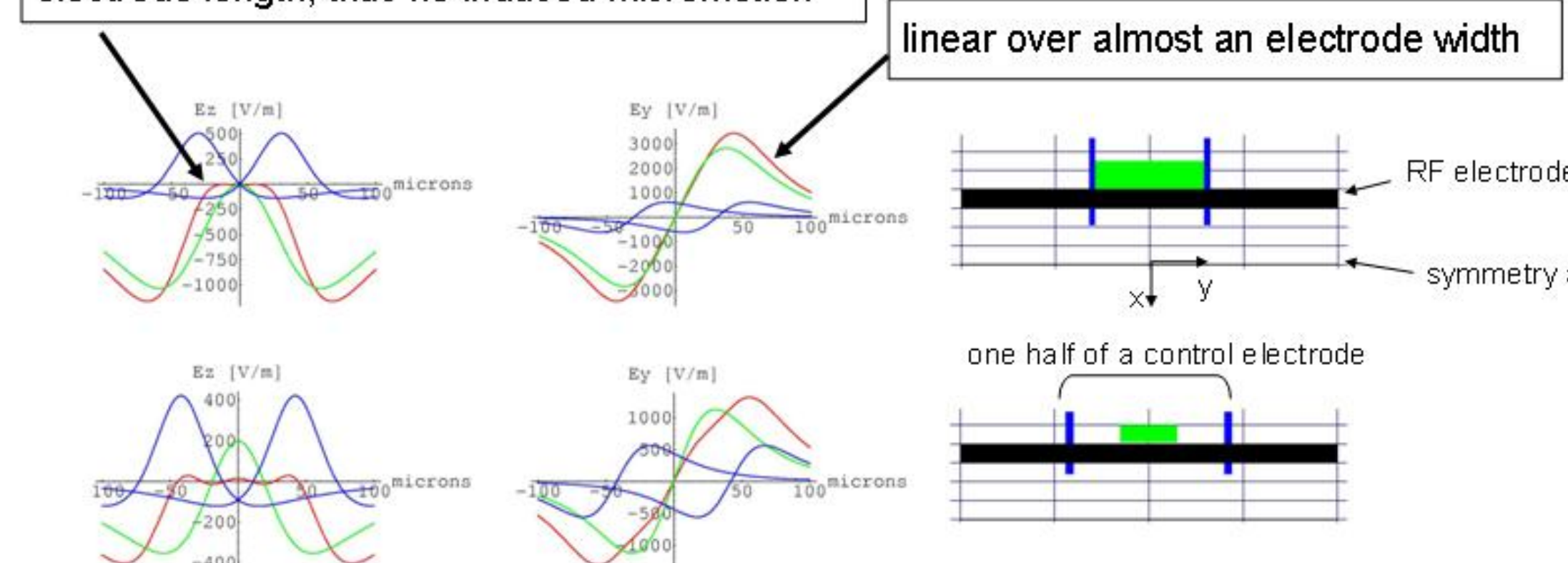
Conclusions:

- trapping depths of symmetric and asymmetric layouts are *identical* for ideal geometry of symmetric four wire trap!
- their t.d.'s are lowered to ~1% of the trap depths of a hyperbolic trap
- frequencies are only slightly higher for the asymmetric layout, and about 0.31 compared to hyperbolic traps
- but: surface traps can be easily manufactured on a much smaller length scale and simpler
- pro&con of asymmetric and symmetric layouts:
symmetric arrangement is simpler to realize technically and preferable in junctions (see below),
principal axes are parallel and perpendicular for symmetric layouts, thus Doppler cooling is hampered

Geometrical compensation of micromotion and nearly parabolic potentials

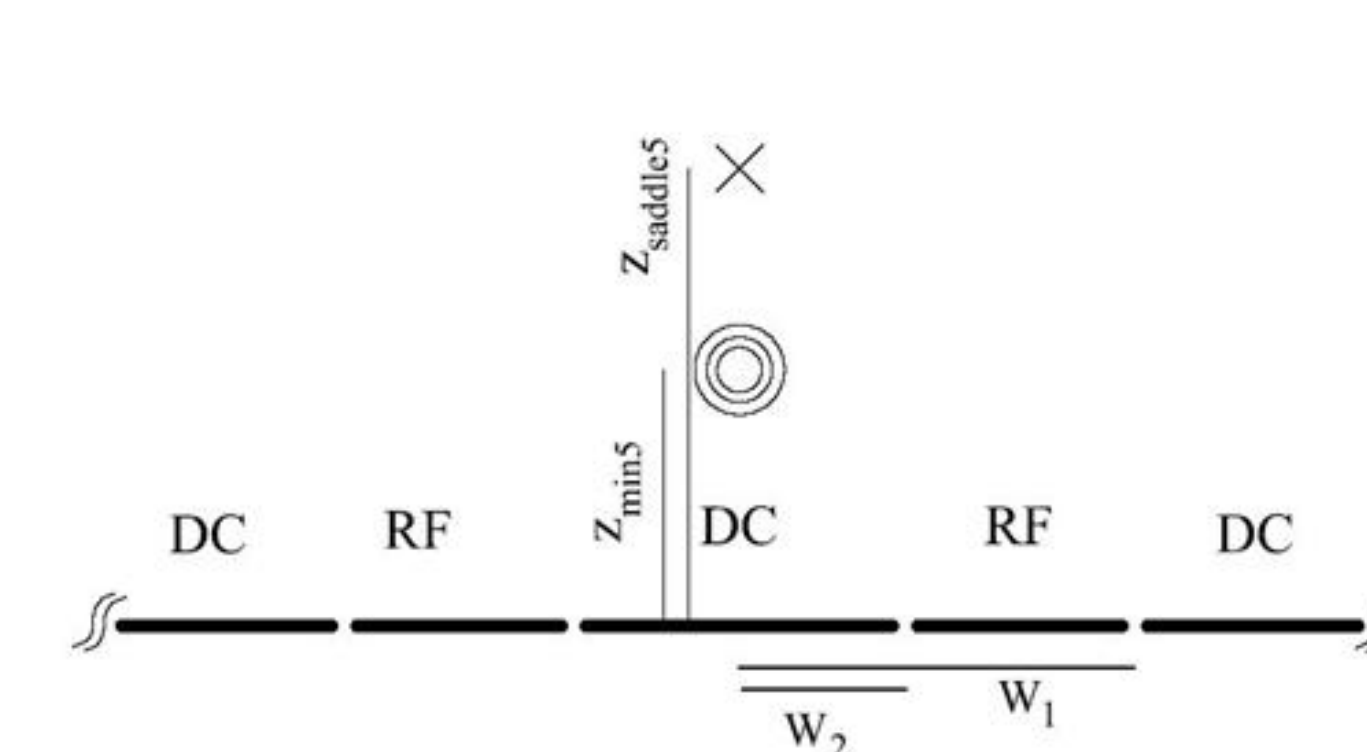
Control electrodes can be designed to have only an E field component in a single planar dimension, e.g. along the trap axis, over a distance longer than the average electrode length!

Ez field (vertical component) is nearly zero over electrode length, thus no induced micromotion



the colors blue and green refer to the electrode parts on the right with the same color. Red represents the whole electrode, i.e. the sum of both parts.

'five wire surface traps' /
(symmetric four wire layouts+ground)



$$z_{\min 5} = \sqrt{W_1 W_2} \equiv G(W_1, W_2)$$

$$z_{\text{saddle}5} = (W_1 W_2)^{1/4} \sqrt{W_1 + W_2 + \sqrt{W_1 W_2}},$$

$$t.d.5(W_1, W_2) = \frac{Q^2 V_{RF}^2}{4\pi^2 m \Omega^2} \times \frac{A(W_1, W_2) - G(W_1, W_2)}{A(W_1, W_2)^2 (A(W_1, W_2) + G(W_1, W_2))}$$

$A(), G()$ arithmetic, geometric mean, respectively

$$\text{maximal for } A(W_1, W_2)/G(W_1, W_2) = \frac{1 + \sqrt{5}}{2} = g \quad (1) \quad \text{GOLDEN RATIO}$$

using the last relation yields $t.d.\text{max}5 \equiv t.d.4$

optimum geometry from (1) $\bar{W}_{1/2} = W_{1/2}/z_{\min 5} = g \pm \sqrt{g}$

$$\omega_z \equiv \omega_x = \frac{QV_{RF}}{\sqrt{2}m\Omega z_{\min 5}^2} \cdot \frac{2}{\pi} \sqrt{\frac{1 + \sqrt{5}}{7 + 3\sqrt{5}}}$$

$$\equiv \omega_{\text{hyp}} \sim 0.309 \text{ reduction factor}$$

Crossing points at junctions (nodal points)

networking requires implementation of junctions

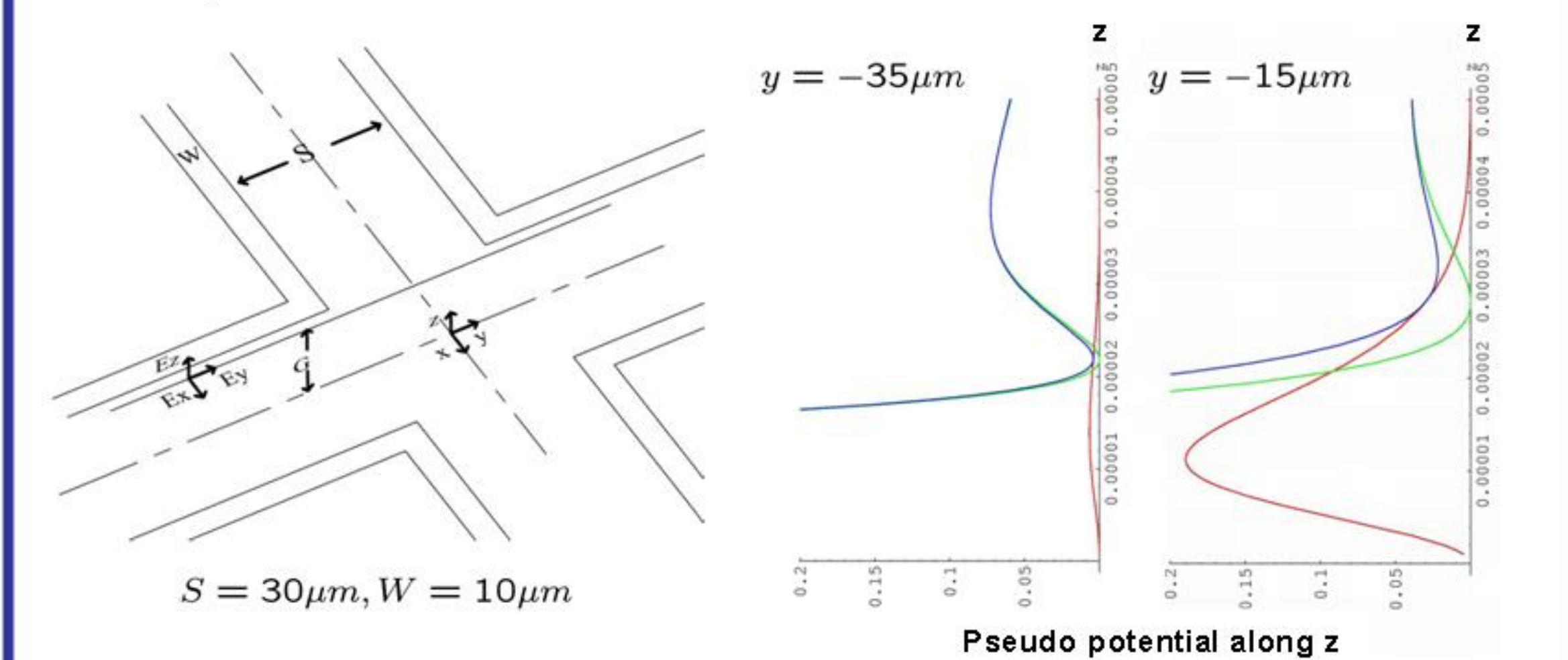
Basic principle: transform from a 2D dynamic 1D static trapping (outside of a junction) to a 1D dynamic 2D static potential (center of junction)

$$\text{Pseudo potential}(x=0) \propto (E_y^2(y, z) + E_z^2(y, z))$$

$$E_y^2 = (E_{y1} + E_{y2})^2$$

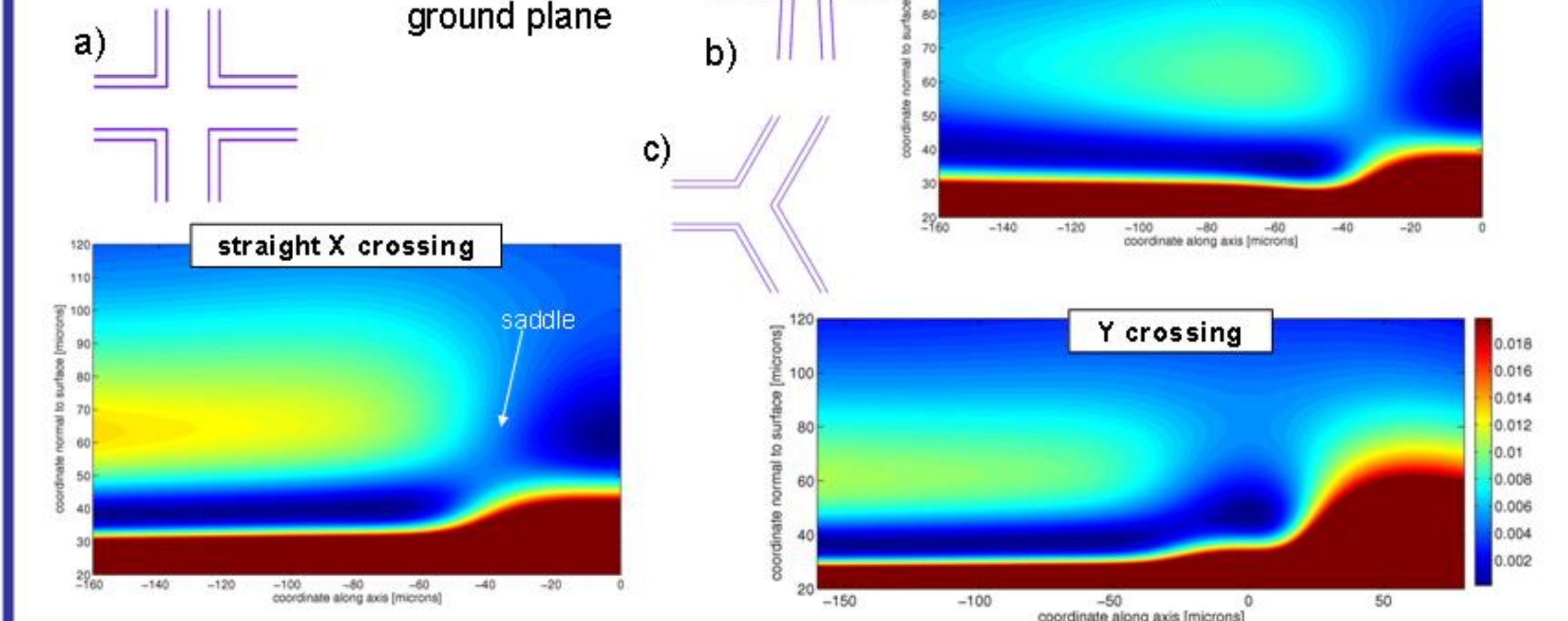
$$E_z^2 = (E_{z1} + E_{z2})^2$$

problems are RF bumps resulting from edge singularities in Ey!
Anti-binding can result



analyzing diverse types of nodal points

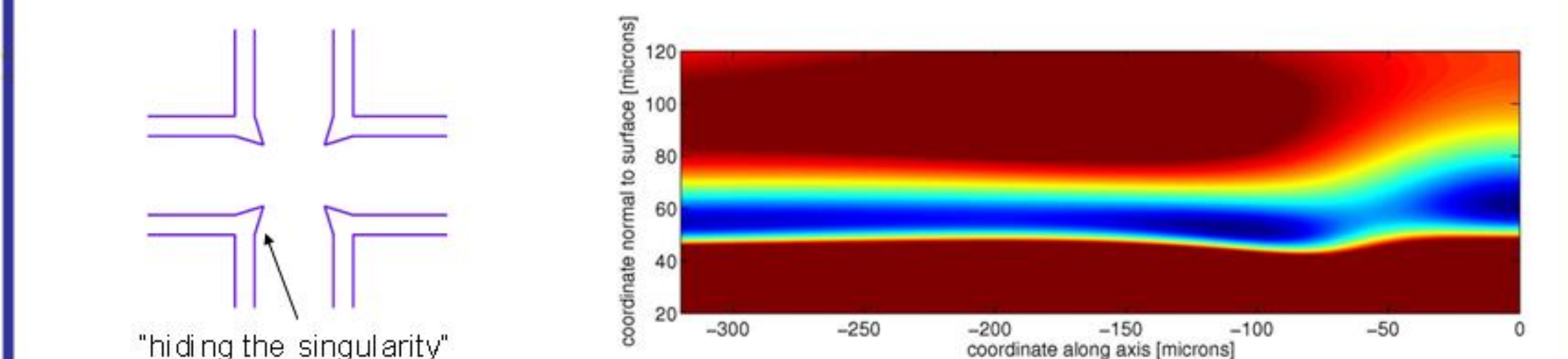
shown: yz contour plots of pseudo-potential for 3 different geometries
sketches: RF rails embedded in a ground plane



Y junctions are preferable and exhibit a smoother behaviour at junctions!

2) another possibility is to use **compensation structures**

modified electrode shapes near the cross help to smooth characteristics near nodal points (reduce effect of singularities)

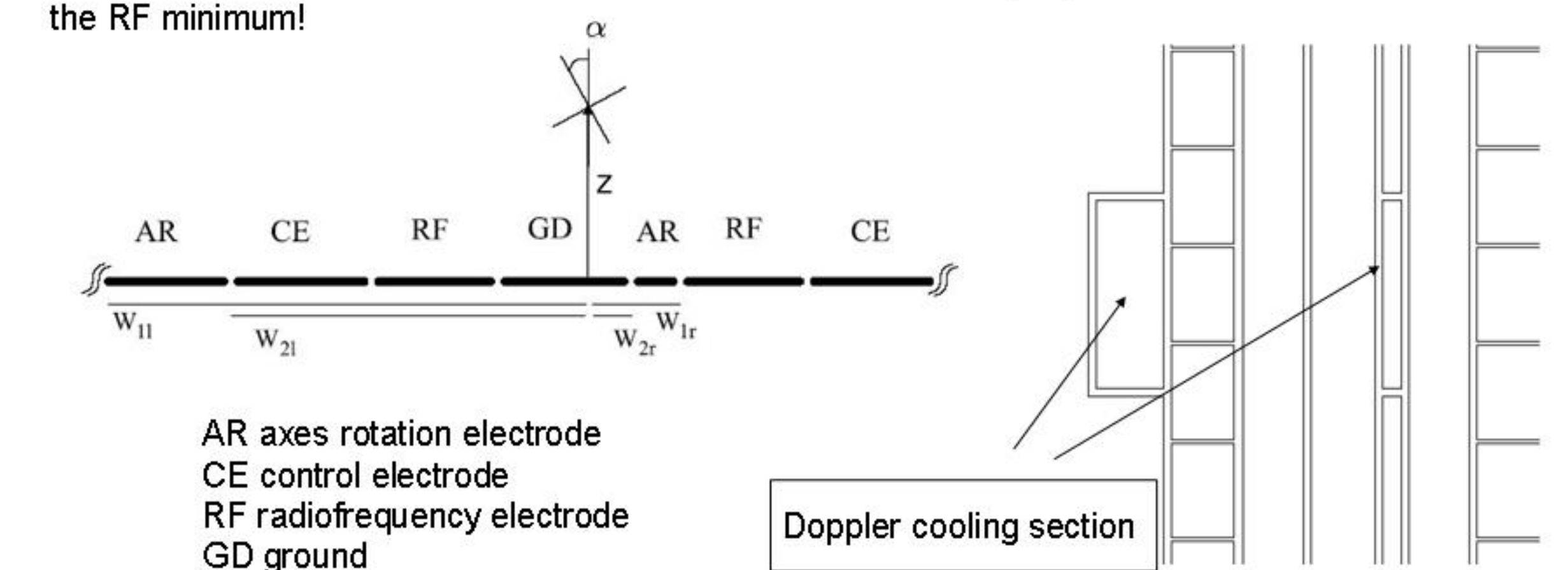


Doppler cooling in 'symmetric four wire traps' (5 wire traps)

From analytical theory follows

$$4z(W_1 + W_2) \tan \alpha - (W_1 + W_2)^2 + (W_1 - W_2)^2/4 + 4z^2 = 0$$

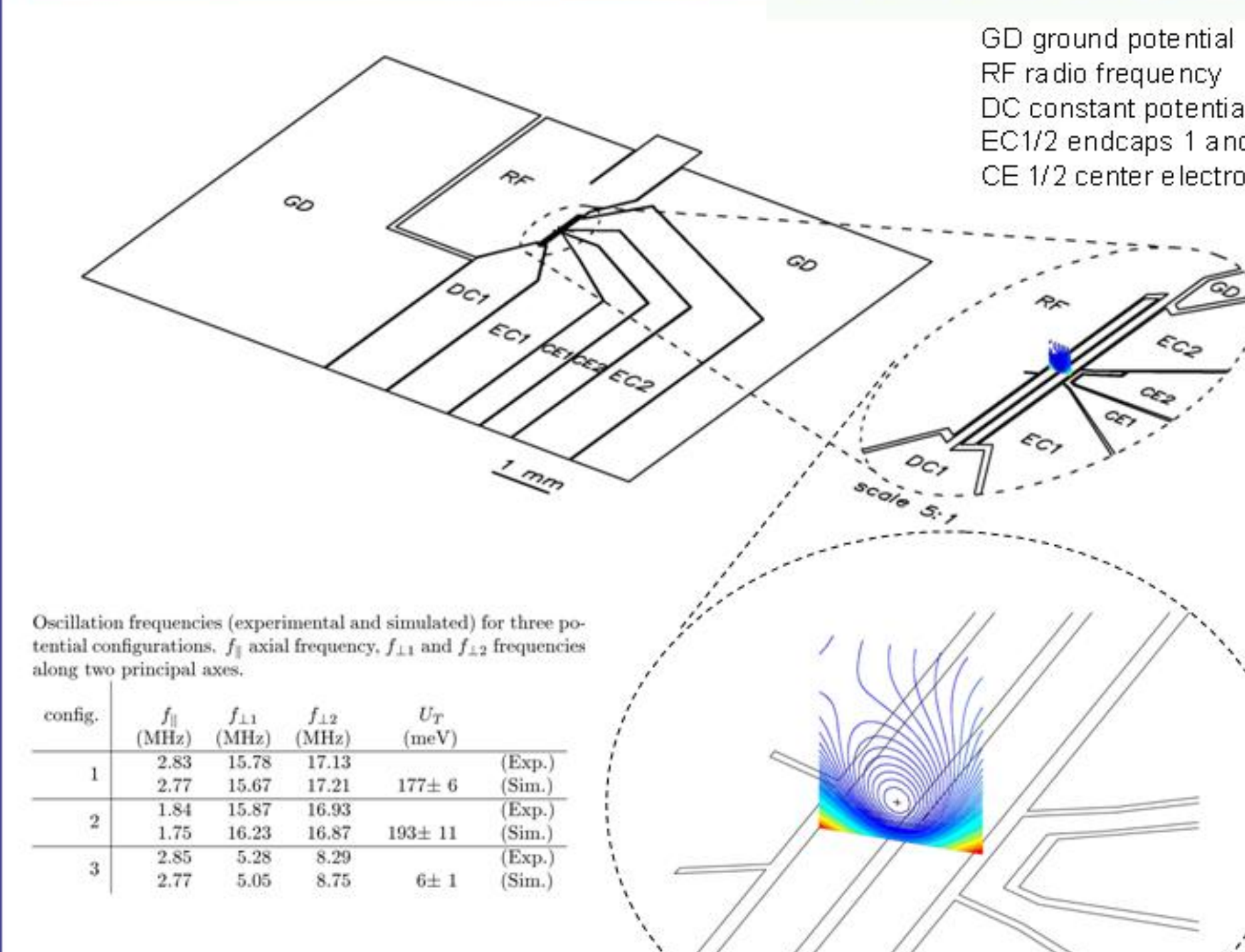
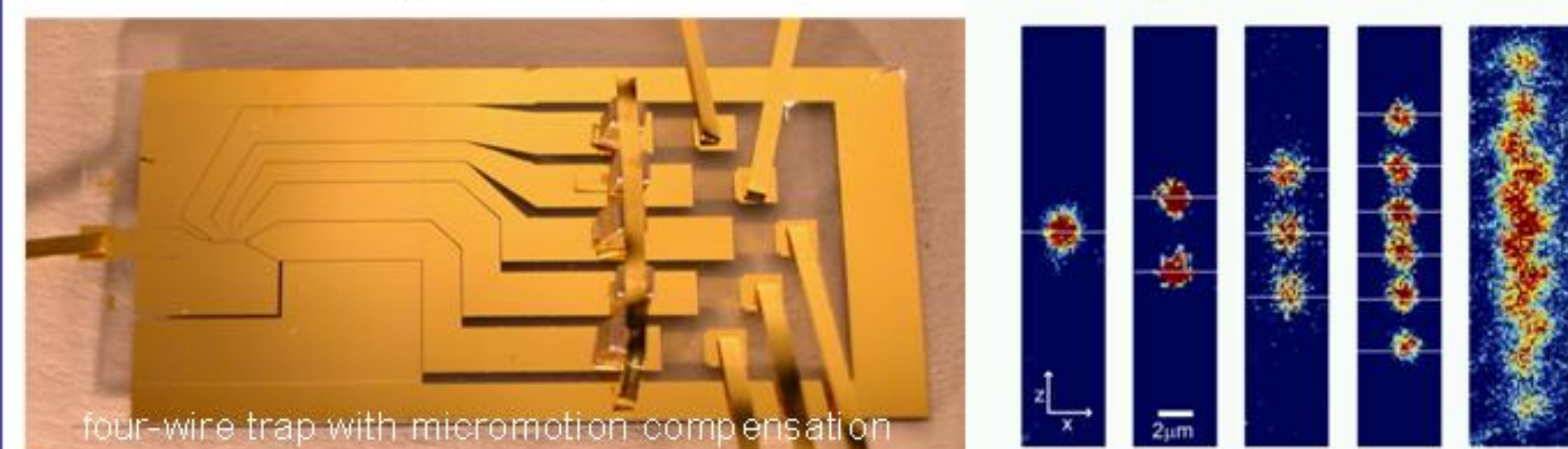
thus we can find locations and widths of 2 additional electrodes (AR) to force a rotated saddle at the RF minimum!



Progress in surface trap layouts for QIP at NIST

(*arXiv quant-ph* 0601173, 2006)

reliable trapping of ions, storing with low heating rate



for similar approaches of other groups see C.E. Pearson *et al.* *arXiv quant-ph* 0511018, (2006)
D. Stick *et al.* *Nature Physics* **2**, 36 (2006)